

Electrostatic lens sizing

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Revision C

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Revision C: correction of bugs and addition of some details

1. Goal, presentation and notations used

The goal of this presentation is to give some additional information about the sizing of electrostatic lenses (i.e. focal length), used to focus particles beams.

The main need of the author was to rapidly size simple lenses, but the different documents found from Internet did not enter in details and did not permit to size the lenses according to the different parameters (voltages and configuration) in a relative simple way, without simulations.

This presentation relies on the Multiplasma simulator program version 1.12 (not public) developed by the author and used for the simulation of the lenses (among other things).

So this paper will permit to size an electrostatic lens by giving an order of magnitude of the focal length.

Notations

- the simple product is indicated with « * » or « x » or « . » or is not indicated if there is no ambiguity,
- the powers of ten are indicated with E_x or 10^x (for example 10^{-7} or E-7),
- the other powers are noted \wedge (for example x^2 for x^2) or “SQR(x)” for x^2 or SQRT(x) for $x^{0.5}$ (square root).
- “<<” for “very inferior” and “>>” for “very superior”
- $|x|$, absolute value of x

The author uses SI units or sub-multiple: the unit of length taken here is the “mm” which is more convenient than the “m” for lenses.

Notes:

- the resulting focal length depends on SQR(U_a/U_I), not on the absolute value of U_a or U_I (see §2 for meaning of U_a and U_I),
- breakdown problem between electrodes is not addressed.

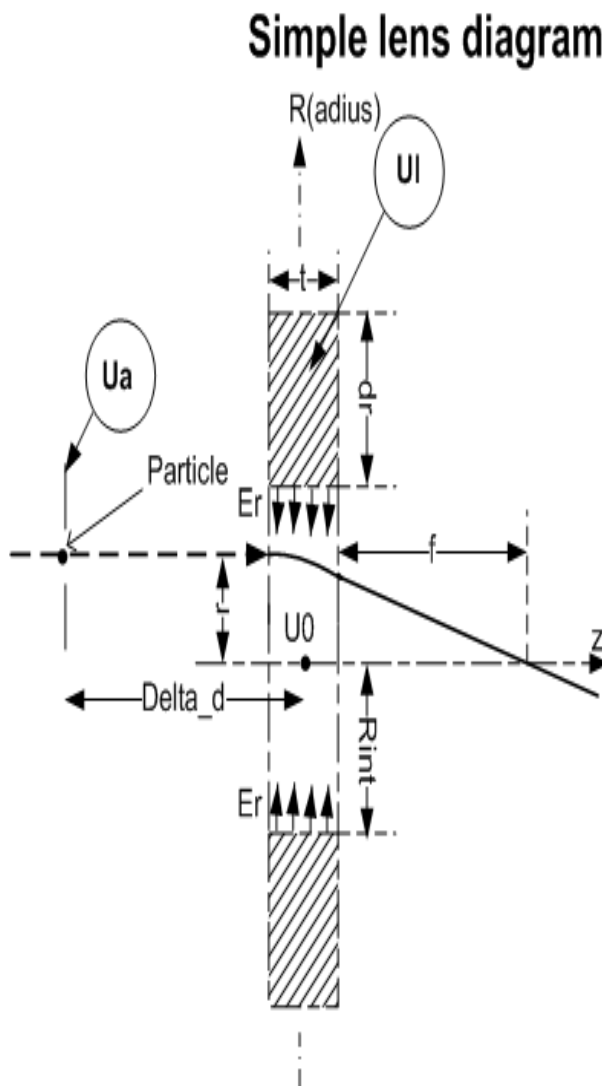
Hypothesis:

- It is supposed that the initial particle trajectory is horizontal. However due to the not nil beam emittance, the beam is normally more or less divergent, so with a non nil radial speed. This will not affect the calculated focal length.
- It is supposed that the space charge can be neglected. Note that if the space charge is large, its effect is predominant and the lens is without any use. The maximum “divergent” space charge radial electric field E_c (V/m) can be determined for a current I (A) of particles beam of radius R (m) and speed V (m/s), by the formula:

$$E_c = I / (2 \cdot \pi \cdot \epsilon_0 \cdot V \cdot R)$$
with $\epsilon_0 = 8.85 \cdot 10^{-12}$ F/m
This value of E_c must be very inferior to the “convergent” radial electric field E_r generated by the lens. A very rough estimation would be based on a U_0 equal to $U_I/2$. So the rough mean radial E_r along the radius would be equal to $U_I / (2 \cdot R \cdot \ln 2)$.
If $E_c \ll E_r$ the lens can be of some use.

2. Description of the simple lens

The goal of the lens is to focus a beam. Due to the interest of the author, it will be supposed that the particles are ions with one charge (D^+ or T^+). But it could be electrons with opposed voltages. Below the diagram shows the different notations. The simple lens is in fact a single electrode, in form of washer. This type of lens can also be called "Wehnelt". Physically, the positive voltage of the lens pushes the ions towards the center (through the radial electric field E_r).



Ul: voltage (V) applied on the lens

Ua: acceleration voltage (V) applied to the particle.

It is equivalent to the energy in eV at which the particle (supposed with a single charge) is sent.

t: thickness (mm)

dr: exterior radius – interior radius (mm)

Rint: interior radius (mm)

f: focal length, in absence of space charge (to calculate) (mm)

r: radius at which the particle penetrates the lens (mm)

Delta_d: distance between the injected particles and the lens center (mm)

Er: radial electric field (V/mm)

U0: induced potential at the lens center (to calculate)

Notes:

* if the particle is an electron, U_a is positive and U_l is negative, with $|U_l| \ll |U_a|$

* if the particle is an ion, U_a is negative and U_l is positive, with $|U_l| \ll |U_a|$

3. Simple lens sizing

We are going, first, to determine a simplified theoretical formula for the focal length, with the first following hypothesis:

- The lens has a cylindrical symmetry, so the azimuthal behavior is not taken into account.
- The trajectory of the particle is initially horizontal (along and above the z axis),
- The distance “r” to the axis is very small compared to the lens radius (Gaussian hypothesis). r is positive by agreement, the axis R carrying “r” being upwards. With this Gaussian hypothesis, the focal length will not depend on “r”. In the reality it is worth, roughly, on the first half of the lens radius. Beyond, the beam converges with different smaller focus lengths (causing aberrations).

It can be demonstrated that the trajectory of the ion obeys to:

$$d^2r/d^2z = -q/(m \cdot \text{SQR}(vz)) \cdot (\delta V/\delta r - \delta V/\delta z \cdot (dr/dz))$$

With r the radius, z the axial distance along the z axis, q the charge in Coulomb, m the mass in kg, vz, the axial speed, V the potential (which depends on r and z).

For details, look at the reference [1], page 15.

Here it will be done the hypothesis that dr/dz (i.e. the angle of the trajectory) will be always very small so as to neglect the second term. This supposes that U_l/U_a will be small.

So the expression can be simplified in: $d^2r/d^2z = -q/(m \cdot \text{SQR}(vz)) \cdot (\delta V/\delta r)$.

Note that $\delta V/\delta r = -E_r$ (the radial electric field)

In addition, it can be demonstrated (divergence theorem) that:

$$\delta V/\delta r = 1/2 \cdot (\delta E_z/\delta z) \cdot r$$

With E_z the electric field along z.

For details, look at the reference [1], page 17.

Moreover, $m \cdot \text{SQR}(vz) = 2 \cdot E_k$ (E_k being the kinetic energy),

If E_k is expressed in eV (and not in Joule), it follows:

$$m \cdot \text{SQR}(vz)/q = 2 \cdot E_{k_{eV}} \text{ (expressed in eV) or } q/(m \cdot \text{SQR}(vz)) = 1/(2 \cdot E_{k_{eV}})$$

In what follows $E_{k_{eV}}$ will be replaced by U_a (in V).

Note: normally E_k (in J) = $q \cdot U_a$ (V) but with a one charge ion, E_k (in eV) = U_a (numerically), U_a being the acceleration voltage of the particle.

So will all these considerations:

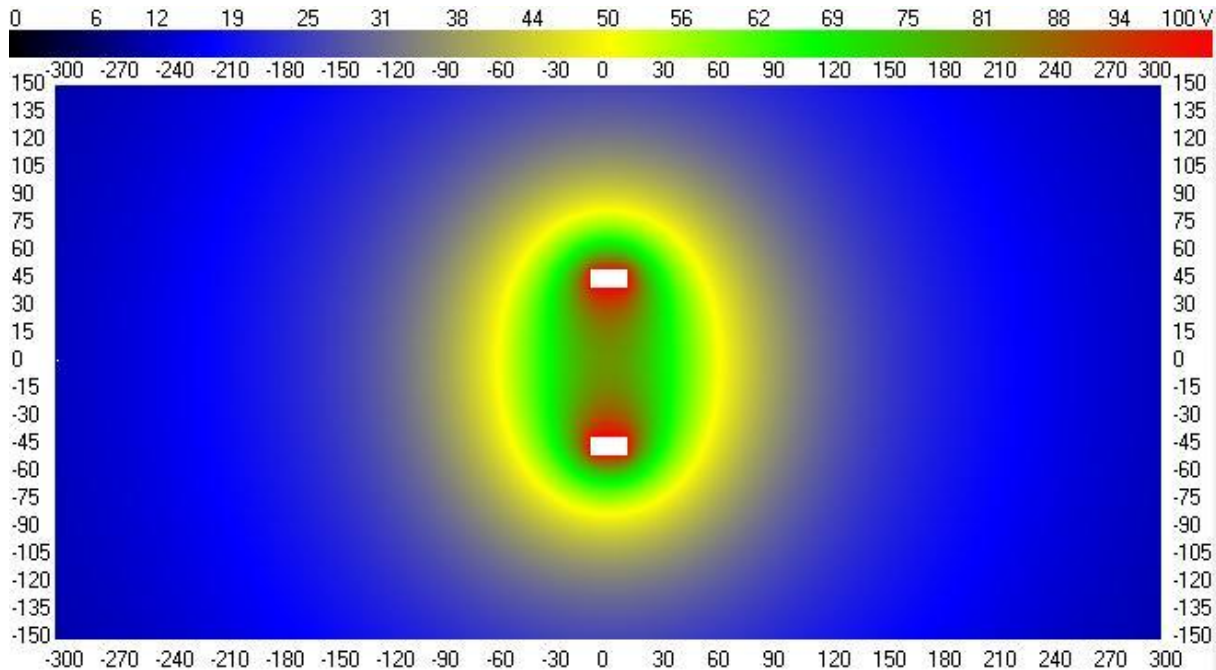
$$d^2r/d^2z = -1/2 \cdot 1/2 \cdot (\delta E_z/\delta z) \cdot r / U_a = -1/4 \cdot (\delta E_z/\delta z) \cdot r / U_a$$

By integration along z, we have $dr/dz = -E_z \cdot r / (4 \cdot U_a)$, equal to the angle α if dr/dz is weak ($\tan(\alpha) \approx \alpha$). With α , it can be determined f (focal length) = $|r/\alpha|$

Moreover, E_z (called E_{z1}) is negative before the lens, and E_z (called E_{z2}) is positive after the lens, with normally $E_{z2} = -E_{z1}$

←---- E_{z1} ---- Lens at + U_0 (V) ---- E_{z2} -----→

Look at the electric field of the lens below (in term of induced potential in V). It appears clearly that, circulating from left to right along z, E_z changes of sign at the center of lens. Note that the electric field tends to become as a spherical electric field when the radius increases.



It can be observed that, from left to right on the axis, the induced potential increases up to U_0 at the center of the lens and then decreases from U_0 .

From all these elements, it can be shown (reference [2] page 116) that $f = 4 \cdot \text{SQR}(U_a) / (E_z \cdot U_l)$

Note: as E_z is proportional to U_l , it follows that f is proportional to $\text{SQR}(U_a/U_l)$.

E_z is not known but it can be determined experimentally. Moreover for a given lens, $U_0 = K \cdot U_l$ (K determined experimentally according to t , R_{int} and dr).

About E_z , experimentally, at $D_1 = 2 \cdot (R_{int} + 2/3 \cdot dr) + 1/2 \cdot e$, the induced voltage is relatively constant around $0.416 U_0$ for the considered lenses. This gives a reference $E_{z0} = 0.584 \cdot U_0 / D_1$.

The real E_z undergone depends mostly on the distance Δ_d . It is equal to $E_z = K_2 \cdot E_{z0}$, It remains to determine the constant K_2 experimentally (according to Δ_d , R_{int} and dr). Note that it is implicitly supposed that the electric field after the lens center is symmetrical of this one (which is wrong but necessary to remain simple).

Finally, f (mm) = $4 \cdot \text{SQR}(U_a) \cdot D_1 \cdot \text{SQR}(K) / (K_2 \cdot 0.584 \cdot U_0 / K \cdot U_0 / K)$, or

$$f$$
 (mm) = $4 \cdot \text{SQR}(U_a/U_0) \cdot D_1 \cdot \text{SQR}(K) / (K_2 \cdot 0.584)$

In the program the constant $4 \cdot \text{SQR}(K) / (K_2 \cdot 0.584)$ is replaced by the constant $3.36 \cdot K_1$, so f (mm) = $3.36 \cdot K_1 \cdot \text{SQR}(U_a/U_0) \cdot D_1$, but it comes to the same, as the K_1 is determined to match, at best, the simulations (supposed to exactly correspond to the reality) with the calculated f .

The complete calculation with the experimental formulas for U_0 and K_1 is given in the Pascal (Delphi 6) procedure in Appendix.

So given U_l , U_a , R_{int} , d_r , t and Δ_d , it will be given the probable focal length (f) with an estimated dispersion (compared to simulations) from 1/2 to 2 times the result.

An example is given in the program in Appendix.

The limits of validation are the following:

- Δ_d between 1 and 10 times the lens exterior diameter. Note that the influence of the acceleration electrode at the voltage U_a is not taken into account (neglected),
- d_r inferior or equal to R_{int}
- e inferior or equal to $2 \cdot R_{int}$
- $U_l \leq 0.4 \cdot U_a$

Physical explanation of the proportionality of f with $\text{SQR}(U_a/U_l)$

Note 1: in this chapter, “K” means “any positive value of any dimension” (this to avoid useless developments). Moreover, “K” is not submitted to algebraic operations.

Note 2: ions are only considered in this chapter, this to always have K positive,

Note 3: the sign “~” is worth for “proportional”

Intuitively, it seemed natural, for the author, to think that the focal length f be proportional to U_a/U_l for the following (bad) reason:

E_r (the radial electric field, convergent so negative) is proportional to $U_l - U_0$ and finally to U_l (U_0 being proportional to U_l).

So the radial speed $v_r = K \cdot E_r \cdot cd$ with “cd” for “crossing duration”.

$cd = t/v_z$ (t : thickness and v_z : axial speed along z). Hence $v_r \sim U_l/v_z$.

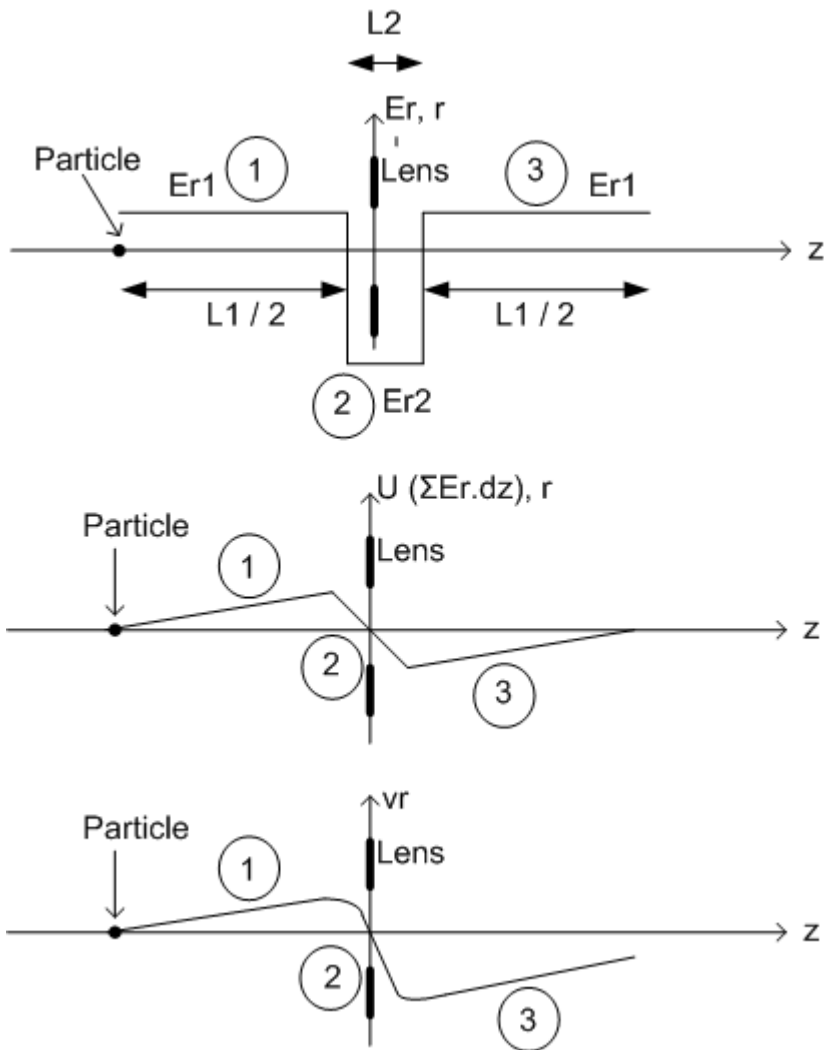
The angle $\alpha = v_r/v_z \sim U_l/\text{SQR}(v_z)$ or U_l/U_a (because $v_z \sim \text{SQRT}(U_a)$)

And finally $f = |r/\alpha| \sim U_a/U_l$

But it is not true, because f is indeed proportional to $1/\text{SQR}(U_a/U_l)$.

Looking the evolution of the electric field along the trajectory, it can be noted that for the most part of the trajectory the electric field is low and positive (divergent), whereas at the lens level and a bit around it, the electric field is high and negative (convergent).

In the diagram of the next page, it is shown the rough and simplified shapes of the electric field E_r , its integration along the trajectory $\sum E_r \cdot dz$ and the radial speed v_r .



It can be noted that $\sum E_{r1} \cdot dz = -\sum E_{r2} \cdot dz$ so $E_{r1} \cdot L1 = -E_{r2} \cdot L2$ or $E_{r2} = -E_{r1} \cdot L1/L2$ (with $L1 \gg L2$).

The parts 1 and 3 ($E_{r1} > 0$) are crossed by the particle at a mean speed v_{z1} along z : $v_{z1} = K \cdot \text{SQRT}(U_a - U_0/2)$, so the time t_1 to cross the parts 1 and 3 is equal to $t_1 = L1/v_{z1}$.

The total positive speed v_{r1} (upwards) got by the particle when it crosses the parts 1 and 3 is equal to $v_{r1} = K \cdot E_{r1} \cdot t_1$ (based on the classical formulas: $v = \text{acceleration} \cdot t$ and $\text{Force} = q \cdot E_r = m \cdot \text{acceleration}$).

So $v_{r1} = K \cdot E_{r1} \cdot L1/v_{z1}$

The part 2 ($E_r < 0$) is crossed by the particle at a mean speed v_{z2} along z :

$v_{z2} = K \cdot \text{SQRT}(U_a - U_0)$, so the time t_2 to cross the part 2 equal to $t_2 = L2/v_{z2}$.

The total negative speed v_{r2} (downwards) got by the particle when it crosses the part 2 is equal to $v_{r2} = K \cdot E_{r2} \cdot t_2$ or $v_{r2} = K \cdot E_{r2} \cdot L2/v_{z2} = -K \cdot (E_{r1} \cdot L1/L2) \cdot L2/v_{z2} = -K \cdot E_{r1} \cdot L1/v_{z2}$

The final speed $v_r = v_{r1} + v_{r2} = K \cdot E_{r1} \cdot L1 \cdot (1/v_{z1} - 1/v_{z2}) = K \cdot E_{r1} \cdot L1 \cdot (v_{z2} - v_{z1}) / (v_{z2} \cdot v_{z1})$

Noting that $U_0 \ll U_a$ and so using $\text{SQRT}(1 - \text{Epsilon}) = 1 - \text{Epsilon}/2$, it comes, neglecting U_0 in front of U_a : $v_r = -K \cdot E_{r1} \cdot L1 \cdot U_0 / (U_a^{1.5})$ which is negative and so convergent (towards the axis).

The angle α of the particle trajectory is equal to $\alpha = v_r/v_{z_final}$
with $v_{z_final}=K.SQRT(U_a)$.

So $\alpha = -K.Er1.L1.U0/(U_a^{1.5})/(K.SQRT(U_a))=-K.Er1.L1.U0/SQR(U_a)$

As $Er1=K.(U1-U0)$, with $U0=K.U1$, it means that $Er1=K.U1$

Finally $\alpha = -K.SQR(U0/U_a)$ or $\alpha = -K.SQR(U1/U_a)$

So f (focal length)=|r/ α | is proportional to $SQR(U_a/U1)$ or $SQR(U_a/U0)$, as expected.

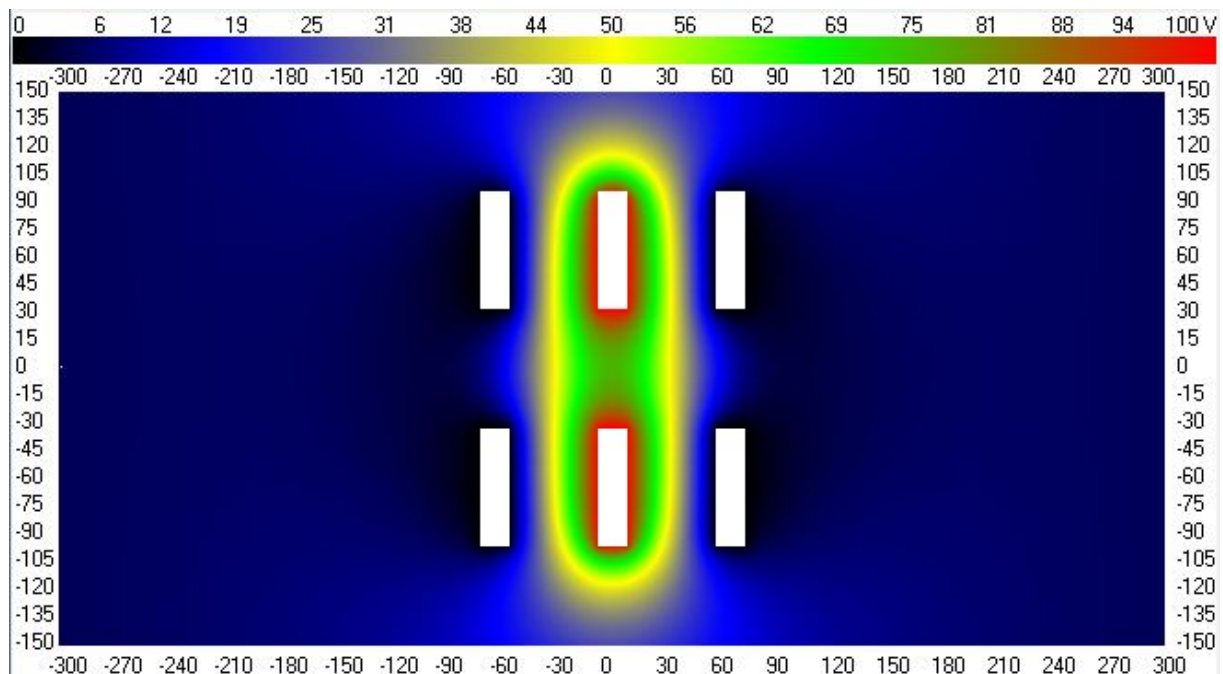
4. About lenses with 0V symmetrical electrodes

It can be simpler to force the electric field with two electrodes at 0 V, symmetrically disposed. The number of parameters being high, it will not be proposed a sizing, but two examples (type 1 and type 2).

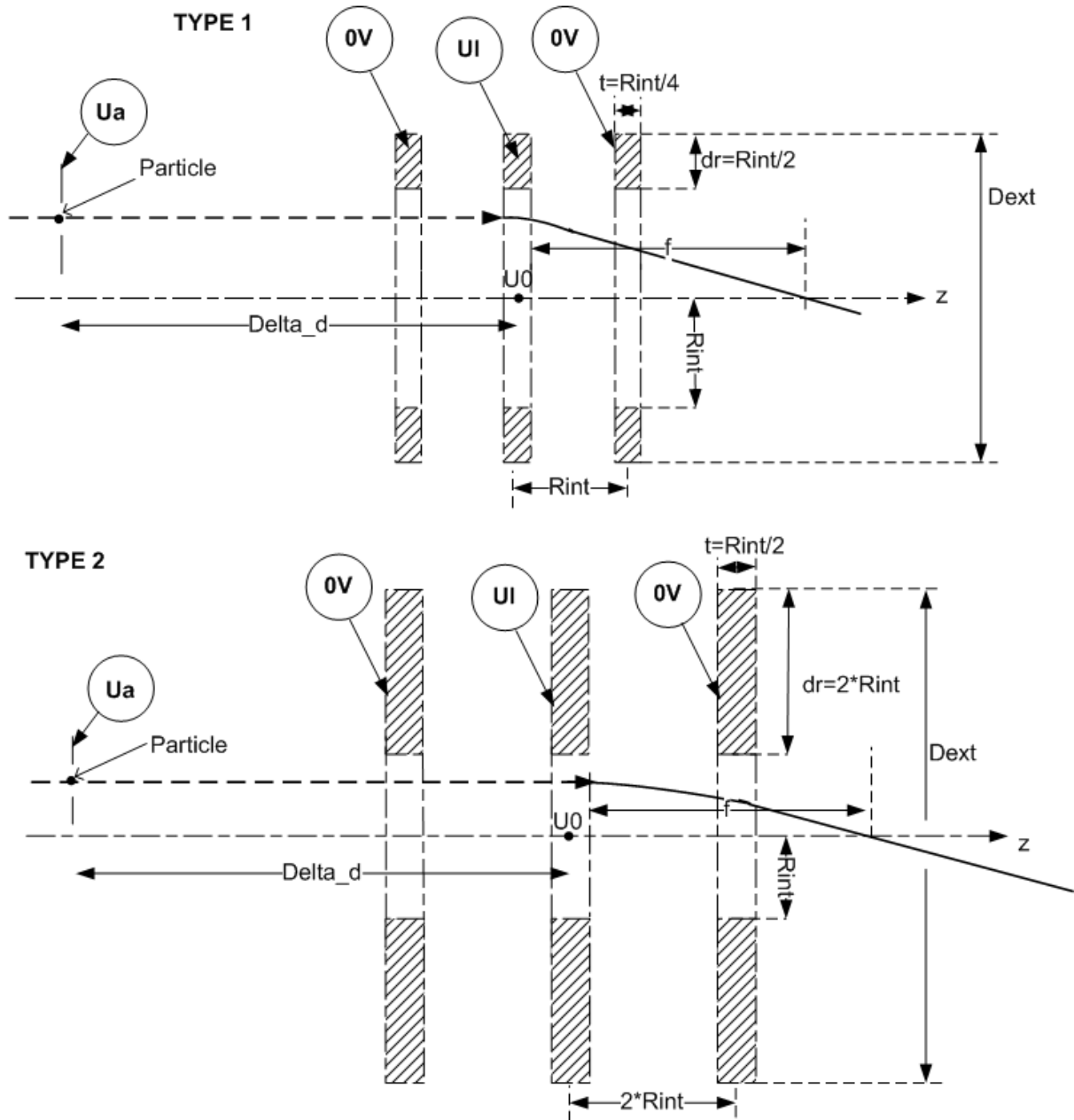
The big advantage of that type of lens is that it does not depend much on the distance between the injected particles position and the lens center (Δ_d), contrary to the single lens. It is almost an isolated system

In the following page, it will be found the diagram of these two examples of lens.

Bellow is an example of electric field for such lenses (in term of induced potential in V). It is obvious, compared to the previous one for a single electrode that the electric field is well limited by the two symmetrical electrodes at 0 V.



Two examples of lenses with symmetrical 0 V electrodes



UI: voltage (V) applied on the lens

Ua: acceleration voltage (V) applied to the particle.

It is equivalent to the energy in eV at which the particle (supposed with a single charge) is sent.

t: thickness (mm)

dr: exterior radius – interior radius (mm)

Rint: interior radius (mm)

Dext: lens external diameter (mm)

f: focal length, in absence of space charge (to calculate) (mm)

Delta_d: distance between the injected particles and the lens center (mm)

U0: induced potential at the lens center (V)

Notes:

* if the particle is an electron, Ua is positive and UI is negative, with $|UI| \ll |Ua|$

* if the particle is an ion, Ua is negative and UI is positive, with $|UI| \ll |Ua|$

The focal length f_1 for the type 1 is given experimentally by the formula:

$$f_1 \text{ (mm)} = 147 * (R_{int}/8) * ((\Delta_d/(3*D_{ext}))^{0.1}) * \text{SQR}(U_a/U_I)$$

As it can be seen, due to the small 0.1 power factor, the influence of Δ_d is weak (but still monotone).

The focal length f_2 for the type 2 is given experimentally by the formula:

$$f_2 \text{ (mm)} = 40 * (R_{int}/4) * \text{SQR}(U_a/U_I)$$

As it can be seen, the influence of Δ_d is not taken into account because it is weak (< to 15%) and, overall, not monotone.

5. Conclusion

- The focal length estimate (f) of the single lens (diagram in page 3) is proposed in Appendix.
- The focal lengths of the type 1 and type 2 lenses with 0 V symmetrical electrodes (diagram in page 9) are proposed just above.

The dispersion of the result (compared to a simulation which has also a certain margin of error) is estimated to be between 1/2 and 2 times the result.

6. References

[1] « Etude théorique et expérimentale de la focalisation des ions afin d'améliorer la brillance du faisceau ionique par suppression des causes d'aberrations » by Jean Faure

[2] « Sur une nouvelle méthode de focalisation des faisceaux d'ions rapides. Application à la spectrographie de masse » by Louis Cartan

APPENDIX

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PROCEDURE SIMPLE_LENS_SIZING_FOR_IONS;
VAR t:SINGLE;{"t" for "thickness of the lens" in mm}
VAR U1:SINGLE;{voltage (in V) applied on the lens (positive for ions and negative for electrons)}
VAR Ua:SINGLE;{accelerator voltage (in V) applied on the particles (negative for ions and positive for electrons)}
VAR Rint:SINGLE;{interior radius in mm}
VAR Dext:SINGLE;{exterior diameter in mm}
VAR dR:SINGLE;{exterior radius - interior radius in mm}
VAR U0:SINGLE; {induced potential at the lens center (to calculate)}
VAR Delta_U:SINGLE;{|U1-U0| in V}
VAR D1:SINGLE;{standard distance compared to the lens center in mm}
VAR Delta_d:SINGLE;{distance between the injected particles and the lens center in mm}
VAR K1:SINGLE;{coefficient to take into account Delta_d, for the Ez estimation}
VAR f:SINGLE;{focal length to calculate in mm}
BEGIN
  {set of hypothesis, for ions}
  Delta_d:=99;
  U1:=40000;
  Ua:=-100000;
  dR:=2;
  Rint:=9;
  t:=4;

  {not permitted}
  IF (U1<0) OR (Ua>0) THEN EXIT;
  IF ABS(U1)>ABS(Ua) THEN EXIT;

  {calculation}
  Delta_U:=U1*0.2479*Power(2.25*t/Rint,-0.2863-(t/Rint)*0.4657)*Power(4.5*dR/Rint,-0.2852-(dR/Rint)*0.0891);
  U0:=U1-Delta_U;
  Dext:=2*(Rint+dR);
  D1:=2*(Rint+2/3*dR)+t/2;
  IF Delta_d<=3*Dext THEN K1:=Power(Delta_d/(3*Dext),0.6-(Delta_d-3*Dext)/(6*Dext))
  ELSE K1:=Power(Delta_d/(3*Dext),0.34);
  f:=3.36*K1*D1*SQR(Ua/U0);

  {display of f}
  //For the previous set of hypothesis, the result is f=966 mm;
  WRITELN('f=',f:4:0,' mm');
END;

```